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**NONCONTACT SENSITIVITY AND COMPLIANCE CALIBRATION METHOD FOR  
CANTILEVER-BASED INSTRUMENTS**

**CROSS-REFERENCE TO RELATED APPLICATIONS**

This application claims priority of U.S. Provisional Application No. 60/272,697, filed on February 28, 2001, the disclosures of which are incorporated fully herein by reference.

**ABSTRACT**

A method for determining physical properties of micromachined cantilevers used in cantilever-based instruments, including atomic force microscopes, molecular force probe instruments and chemical or biological sensing probes. The properties that may be so determined include optical lever sensitivity, cantilever spring constant and cantilever-sample separation. Cantilevers characterized with the method may be used to determine fluid flow rates. The method is based on measurements of cantilever deflection resulting from drag force as the cantilever is moved through fluid. Unlike other methods for determining such physical properties of cantilevers, the method described does not depend on cantilever contact with a well-defined rigid surface. Consequently, the method may be employed in situations where such contact is undesirable or inconvenient. The method has numerous applications, including molecular force measurements, atomic force microscopy and manipulation technology, chemical or biological sensing,

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[lithographic manufacturing, nanometer scale surface profiling] and other aspects of nanotechnology.

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## **BACKGROUND OF THE INVENTION**

The present invention relates generally to methods for determining physical properties of micromachined cantilevers used in cantilever-based instruments, including atomic force microscopes, molecular force probe instruments and chemical or biological sensing probes. It also involves using cantilevers characterized with such methods to measure fluid flow rates.

Calibrating the sensitivity of cantilevers used in cantilever-based instruments is critical for correctly interpreting the results obtained from such instruments. The foundation of this calibration is the determination of the optical sensitivity of the cantilever, the relationship between deflection of the cantilever and movement of the tip

of the cantilever in the z direction. For this purpose, deflection of the cantilever is measured with the optical detection means common in such instruments, a position sensor collecting light reflected off the back of the cantilever. Knowing the optical sensitivity of the cantilever, the spring constant of the cantilever may be readily calculated.

The conventional methods for determining optical lever sensitivity have either been destructive or required that the lever be brought into hard contact with a well-defined rigid surface. Because the distances are typically less than one micron, and the relative positions of the cantilever tip and such a surface difficult to locate, making such contact is far from a trivial proposition. The difficulty of the procedure is enhanced by the fact that slippage of the tip laterally over the surface introduces serious errors.

Even if the conventional methods were easy of execution, there are many instances when it is not desirable or convenient to measure optical lever sensitivity by touching a surface. When the results depend on chemical or biological sensitization of the cantilever tip, or if the tip is particularly sharp, hard contact or any contact before performing the experiment may compromise the results. Similarly compromising may be hard contact when the sample is coated on the surface and is a soft material such as cells. Finally, in the case of chemical or biological sensing probes, there may not be a rigid surface anywhere near the cantilever against which to press.

Two methods for determining optical lever sensitivity not requiring hard contact with a well-defined rigid surface have been proposed, but each has important limitations and is as yet untested. D'Costa and Hoh proposed estimation of optical lever sensitivity by moving the spot across the position sensor a known distance. Because it is not sensitive to actual motion of the cantilever, this method does not account for differences in cantilever geometry or changes in the alignment of the spot on the lever. These issues become even more critical as the length scale of cantilevers shrink. Sader proposes to rely on a plan view of the lever and the measured resonant frequency and quality factor to estimate optical lever sensitivity spring constant.

Although there has been a large amount of work dedicated to the cantilever calibration issue, the precision of the resulting techniques seems to be limited to 10%. In this situation, it is desirable to make use of another method to check for consistency.

## **SUMMARY OF THE INVENTION**

An object of the present invention is to provide a simple method for calibrating the sensitivity of any type of cantilever used in cantilever-based instruments without making contact with any surface.

A second object is to provide a method for the cantilever to approach a sample surface in a gentle and repeatable manner.

Another objective is to provide a method for calibrating the sensitivity of cantilevers that is easily automated.

Another objective is to measure fluid flow rates using micromachined cantilevers.

These and other objects are achieved according to the present invention by (i) measuring the deflection of the cantilever as it moves at a measured velocity through a fluid, (ii) determining the resonant frequency and quality factor of the cantilever by measuring its thermal spectrum and (iii) deriving optical lever sensitivity from combining these measurements.

### **BRIEF DESCRIPTION OF THE DRAWINGS**

FIG. 1: Prior Art showing conventional calculation of optical lever sensitivity.

FIG. 2: resulting from oscillation of the cantilever base (frequency=40Hz).

FIG. 3: Power spectra of cantilever thermal fluctuations

FIG. 4: Hysteresis loops as oscillating cantilever base is separated from sample.

FIG. 5: Dependence of deflection, resonance and quality factor on separation from sample.

FIG. 6: Dependence of Kappa on separation from sample.

FIG. 7: Spring constant values with different methods.

FIG. 8: Fluid flow around cantilevers.



## DESCRIPTION OF THE PREFERRED EMBODIMENTS

The optical sensitivity of the micromachined cantilever, the derivative of the change in cantilever deflection with respect to change in the  $z$  position of the cantilever tip (typically abbreviated as "OLS"), is the foundation for correctly interpreting the results obtained from cantilever-based instruments. Panel A of FIG. 1 depicts the one of the conventional methods for determining OLS. As shown in the panel, the cantilever is pressed into a hard surface (typically freshly cleaved mica) by the instrument (not shown) and moved an arbitrary distance measured by the instrument. Deflection of the cantilever resulting from this change in position is measured with optical detection means commonly employed in such instruments: low coherence light is focused onto the back of the cantilever with an adjustable focus lens and the light reflecting off the cantilever is collected by an adjustable mirror and guided onto position sensor. The position sensor provides a voltage that is proportional to the deflection of the cantilever.

Panel B of FIG. 1 graphs the deflection of the cantilever vs. the  $z$  position of the tip. As shown in Panel B, it is typical to calculate optical sensitivity as the inverse of OLS ("InvOLS"), the derivative of change in the  $z$  position of the cantilever tip with respect to change in cantilever deflection.

Knowing InvOLS permits us to calculate the cantilever spring constant,  $k$ , from the Equipartition of Energy Theorem:

$$\frac{1}{2}k_B T = \frac{1}{2}k \langle A^2 \rangle \quad (1)$$

where  $k_B$  is Boltzmann's constant,  $T$  is the temperature,  $k$  is the cantilever spring constant and  $\langle A^2 \rangle$  is the mean squared cantilever amplitude ( $\langle A^2 \rangle = \text{InvOLS}^2 \cdot \Delta V^2$ , where  $\Delta V$  is cantilever deflection in volts).

As previously noted, it is not always desirable or convenient to determine InvOLS by making hard contact with a well-defined rigid surface as shown in Panel A of FIG. 1. The invention disclosed here permits determination of InvOLS without touching a surface by measuring cantilever deflection resulting from drag force as the cantilever is moved through a fluid (including air).

A cantilever moving through a fluid will be deflected by a viscous drag force. The measured cantilever deflection is converted to a force using  $F_{hyst} = k \cdot \text{InvOLS} \cdot \Delta V$ , where  $\Delta V$  is cantilever deflection in volts as measured by the position sensor. If the cantilever is moving at a speed  $v$  through the fluid, we can characterize the dissipative force, which absent turbulence is equal to the drag force, as  $F_{hyst} = -b_{hyst} v$ , with  $b_{hyst}$  being the damping coefficient. FIG. 2 shows a typical hysteresis loop measured by a 40Hz sinusoidal cycling of the base position of a cantilever while monitoring cantilever deflection with a position sensor.

An independent measurement of the damping coefficient can be made by observing the thermal fluctuations of the cantilever. The simple harmonic oscillator model gives the damping coefficient in terms of the spring constant  $k$ , the resonant frequency  $\omega_0$  and the quality factor  $Q$  as;

$$b_{therm} = \frac{k}{\omega_0 Q} \quad (2)$$

FIG. 3 shows four power spectra of cantilevers in fluid. The two high frequency curves were made in air, one relatively far away from the surface, the other relatively close. The spectrum taken close to the surface had increased damping, yielding a peak with a lower quality factor. The two curves with low peak frequencies were taken in fluid, which caused the resonance to be significantly damped. The fluid is also carried along with the lever as it moves, creating an effective mass that lowers the resonant frequency. The measured spring constants of the four levers (using the method of Hutter and Bechoefer, which requires hard contact with a surface) are virtually the same despite the different environments.

From the data derived from the calculation of such power spectra, and using Equation 2, the damping coefficient,  $b_{therm}$ , may be calculated.

As is well known (see for example, Landau and Lifschitz, *Fluid Mechanics*), damping is a complicated function of the geometry of the damped system, fluid

properties and the amplitude and frequency of the motion. However, for micromachined cantilevers in cantilever-based instruments, we have found experimentally that the thermal and hysteretic damping coefficients are related by  $b_{hyst} = \kappa \cdot b_{therm}$ , where  $\kappa$  is a phenomenological factor that depends on the fluid properties, drive frequency, tip-sample separation and the specific lever geometry. This relationship between the thermal and hysteretic damping coefficients allows us to combine the expressions for viscous drag force and dissipative force in an expression for InvOLS, all in terms of variables measured away from a surface:

$$InvOLS_{hyst} = \frac{\kappa \nu}{\omega_0 Q \Delta V}. \quad (3)$$

Once  $InvOLS_{hyst}$  has been determined, the spring constant can be calculated from the rearranged Equipartition of Energy Theorem:

$$k = \frac{k_B T}{InvOLS^2 \Delta V^2}. \quad (4)$$

FIG. 4 shows a series of hysteresis loops measured at different cantilever tip-sample separations. As the tip approaches the surface (towards the left of the Figure), the damping increases, until, in the last loop, intermittent contact with the surface is made. Panel A of FIG. 5 shows the amplitude of the hysteretic damping as a function of tip-sample separation for a range of cantilever base excitation amplitudes and frequencies. These data were extracted from a number of measurements similar to those shown in Figure 4. Panel B of

FIG. 5 shows the resonant frequency measured with thermal noise as a function of the tip sample separation for a series of similar triangle cantilevers. Panel C of FIG. 5 shows the quality factor  $Q$  measured in a similar fashion as a function of tip-sample separation. The x-axis in all three graphs extends out to 2mm. It is apparent that the three quantities,  $\Delta K$  (Panel A of Figure FIG. 5),  $Q$  (Panel B of FIG. 5) and  $\omega_0$  (Panel C of FIG. 5) have differing dependences on tip-sample separation. This implies that  $\kappa$  is a function of tip-sample separation. FIG. 6 shows  $\kappa$  as a function of tip-sample separation for eleven different amplitudes, frequencies and cantilevers. These measurements imply that  $\kappa$  has a predictable behavior, at least for a particular lever in the amplitude and frequency range tested in this work. The curve in this Figure can then be used in conjunction with Equation 3 to predict *InvOLS* and, via Equation 4, the spring constant of the lever.

FIG. 7 shows five separate calculations of cantilever spring constants for 12 different cantilevers using five different methods, including the hysteretic method disclosed here, as well as the average of all calculations. The variation is pronounced, but each method yields a result generally within 20% of any other.

In the present embodiment, the fluid flow around the cantilever is induced by moving the base of the cantilever. It is also possible, and in some situations desirable, to induce fluid flow around the lever with some other method, such as an external pump or perfusion apparatus. FIG. 8 schematically illustrates such an arrangement, where the fluid flow is controlled by an external apparatus.

The measurement of hysteresis loops while monitoring cantilever deflection required for determining InvOLS in the disclosed method, has an additional benefit. Because the cantilever damping changes as a function of tip-sample separation (see FIG. 4 and the plotted amplitude of hysteresis in Panel A of FIG. 5) observing these hysteresis curves allows the position of the surface to be determined without actually making contact with the sample. The increasing hysteresis as the cantilever approaching the surface allows the surface to be approached to within a few microns without actually making contact. This has utility, for example, in atomic force microscopy where gently bringing the cantilever tip into proximity with the surface is a common challenge.

It is to be noted that if instead of the above situation, the spring constant and damping constant were known but the fluid flow speed was not, the drag measurement would yield a value for the fluid flow rate past the lever via the relationship

$$v = -\frac{b_{hyst}}{k \cdot \Delta V \cdot InvOLS}.$$

This can be used as a probe of fluid flow as illustrated in FIG. 8.

The described embodiments of the invention are only considered to be preferred and illustrative of the inventive concept. The scope of the invention is not to be restricted to such embodiments. Various and numerous other arrangements may be devised by one skilled in the art without departing from the spirit and scope of the invention.